

Chapter III: Electrostatic Fields and Potential

1 – Introduction

The electrostatic potential $V(M)$ associated with the electrostatic field $\vec{E}(M)$ is a scalar function unlike \vec{E} . We will see, in many cases, that the potential will be a convenient intermediary in the calculation of the vector field $\vec{E}(M)$. The potential is physically linked to the concept of potential energy, hence its name.

2 - Circulation of the electrostatic field

2-1 Electrostatic potential

a/ Case of a point load

Consider a point charge $q (>0)$ fixed at P and a point M in space (figure 1):

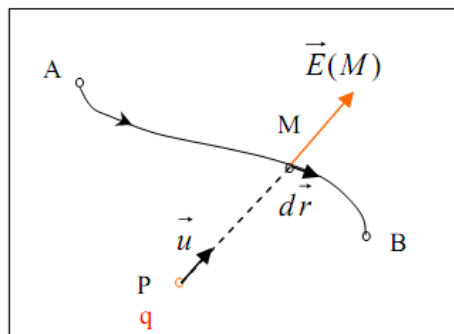


Figure 1

The point charge q fixed at P creates at every point M in space an electrostatic field given by:

$$\vec{E}(M) = \frac{q}{4\pi\epsilon_0} \frac{\overrightarrow{PM}}{\|\overrightarrow{PM}\|^3} = \frac{q}{4\pi\epsilon_0} \frac{\vec{u}}{r^2}$$

$$\text{avec } \vec{u} = \frac{\overrightarrow{PM}}{\|\overrightarrow{PM}\|} : \quad \text{Unit vector directed from P to M.}$$

The elementary circulation DC of the field E corresponding to an elementary displacement $d\vec{r}$ at Point M on the curve AB is:

$$dC = \vec{E} \cdot d\vec{r} = \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \vec{u} \cdot d\vec{r} \quad (1)$$

Or, $d\vec{r} = d(\overrightarrow{PM}) = d(r\vec{u}) = dr\vec{u} + r d\vec{u}$ et $d\vec{r} \cdot \vec{u} = (dr\vec{u} + r d\vec{u}) \cdot \vec{u} = dr + r d\vec{u} \cdot \vec{u}$

Puisque : $d(\vec{u} \cdot \vec{u}) = 2\vec{u} \cdot d\vec{u} = 0$; on a : $d\vec{r} \cdot \vec{u} = dr$

The elementary circulation DC can then be written as:

$$dC = \frac{q}{4\pi\epsilon_0} \frac{dr}{r^2} = -d\left(\frac{q}{4\pi\epsilon_0} \frac{1}{r}\right) \quad (2)$$

Let us then state;

$$dC = \vec{E} \cdot d\vec{r} = -dV(r)$$

V is the electrostatic potential V(M) created by the charge q fixed at M:

$$V(M) = V(r) = \frac{q}{4\pi\epsilon_0} \frac{1}{r} + cste \quad (3)$$

We have just defined a new field, the electrostatic potential; it is a scalar field defined up to a constant. The value of this constant is generally chosen such that the potential is zero when point M is infinitely far from the charge.

$V(r \rightarrow \infty) = 0$. In this case, the constant is zero and the potential can be written as:

$$V(M) = V(r) = \frac{q}{4\pi\epsilon_0} \frac{1}{r} \quad (4)$$

Like the field \vec{E} The potential V is not defined at points P_i : $\vec{E}(P_i)$ et $V(P_i)$ are not defined.

b/ Case of a distribution of n point charges

Let n point charges $q_1, q_2, \dots, q_i, \dots, q_n$ be fixed at points $P_1, P_2, \dots, P_i, \dots, P_n$. Let M be a point in space. (figure 2).

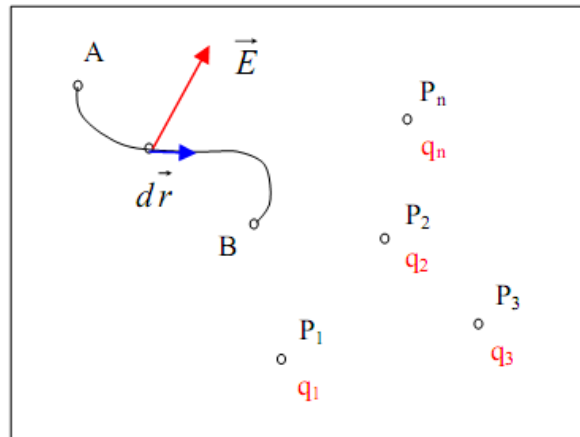


Figure 2

Let's calculate the elementary circulation dC_i of the field \vec{E}_i created by the charge q_i alone:

$$dC_i = \vec{E}_i \cdot d\vec{r} = -dV_i(r)$$

$$\text{Avec } \vec{E}_i(M) = \frac{q_i}{4\pi\epsilon_0} \frac{\overrightarrow{P_iM}}{\|\overrightarrow{P_iM}\|^3} = \frac{q_i}{4\pi\epsilon_0} \frac{\vec{u}_i}{r_i^2} \text{ et } \overrightarrow{P_iM} = r_i \vec{u}_i$$

Thus, the electrostatic potential $V_i(M)$ due to the charge q_i .

$$V_i(M) = \frac{q_i}{4\pi\epsilon_0} \frac{1}{r_i}$$

$$\text{avec: } r_i = \|\overrightarrow{P_iM}\|$$

The potential $V(M)$ due to the set of n charges is the sum of the potentials according to the superposition principle:

$$V(M) = \sum_{i=1}^n V_i(M) = \sum_{i=1}^n \frac{q_i}{4\pi\epsilon_0} \frac{1}{r_i} \quad (5)$$

In this relationship, we have chosen the constant to be zero for each potential V_i created by the charge q_i ; this is only valid if the charges q_i are distributed in a finite volume